

**EUR 4516 e**

COMMISSION OF THE EUROPEAN COMMUNITIES

FOLD

A CALIBRATION PROGRAM FOR THE  
ANALYSIS OF QUASI-ELASTIC NEUTRON  
SCATTERING DATA

by

D.J. WINFIELD

1970



Joint Nuclear Research Center  
Ispra Establishment - Italy

Reactor Physics Department  
Experimental Neutron Physics



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Plots of quasi-elastic peaks calculated by the numerical convolution of the instrumental resolution, which may be of any shape, with Lorentzian functions of various widths are produced by the program. These plots allow a comparison to be made with experimental quasi-elastic peak shapes. The validity of the assumption of a Lorentzian function for the quasi-elastic scattering may thus be tested.

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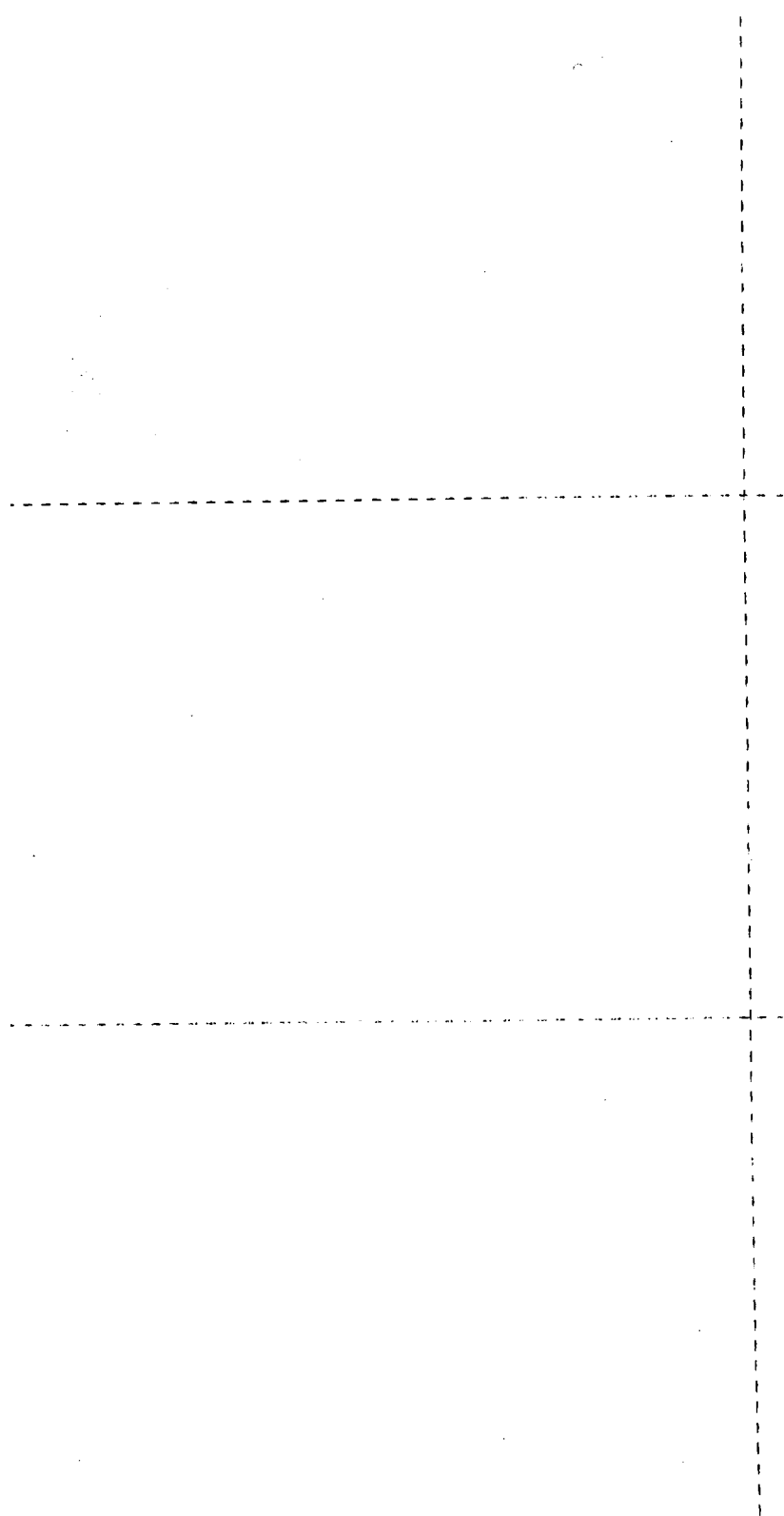
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## **ABSTRACT**

A Fortran IV computer program has been written which gives a means of calibrating the quasi-elastic peaks obtained from the time-of-flight data of the ISPRA-1 Double Chopper Spectrometer in terms of the widths of Lorentzian distributions.

Plots of quasi-elastic peaks calculated by the numerical convolution of the instrumental resolution, which may be of any shape, with Lorentzian functions of various widths are produced by the program. These plots allow a comparison to be made with experimental quasi-elastic peak shapes. The validity of the assumption of a Lorentzian function for the quasi-elastic scattering may thus be tested.

## **KEYWORDS**

FORTTRAN  
COMPUTERS  
CALIBRATION  
ELASTIC SCATTERING  
PEAKS

SPECTROMETERS  
TIME OF FLIGHT METHOD  
CHOPPERS  
LORENTZ LINE SHAPES

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1. Introduction
2. Theory
3. Program Description
4. The Calibration Curve.

References.

Appendix: Program Listing and Typical Output Data.





# 1. Introduction \*)

Diffusive motions of hydrogen containing molecules or hydrogen atoms in samples such as hydrogenous liquids or metallic hydrides may be conveniently studied by the technique of inelastic neutron scattering. These motions generally result in a "quasi-elastic" broadening of the infinitely sharp elastic peak which would be observed in the case of a solid, provided that the thermal motions were zero and the instrumental resolution was perfect. From this broadening, information may be gained about the exact nature of the diffusive motions in the particular sample of interest.

Experimentally however, finite instrumental resolution always imposes a limit on the energy transfers which are observable. With present chopper time of flight spectrometers the smallest detectable energy transfers are about 0.05 meV. The largest quasi-elastic broadenings of interest are normally around 5 meV so that this energy transfer and the above lower limit cover a time scale from  $1.3 \times 10^{-13}$  sec to  $1.3 \times 10^{-11}$  sec. As typical incident energies are about 5 meV with resolutions  $\sim 0.5$  meV, it is therefore always necessary to correct the experimental quasi-elastic peak widths for instrumental resolution in order to extract the true widths of the quasi-elastic peaks, resulting from the broadening process alone.

For a given instrumental resolution function and assuming a scattering law [1]

$$S(\underline{k}, \omega) = \frac{\Delta E_L / 2}{(\Delta E_L / 2)^2 + (\hbar \omega)^2} \quad [1]$$

of Lorentzian form, the program calculates and plots theoretical quasi-elastic peak shapes for a series of

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\*) Manuscript received on 16 June 1970

widths <sup>(\*)</sup> $\Delta E_L$  of this Lorentzian function. Constants have been omitted from equation [1] as the absolute value of the scattering law is not required.

Depending upon the theoretical model used to describe the sample under observation, the Lorentzian width  $\Delta E_L$  is usually a function of the momentum transfer and it is this  $\Delta E_L(\underline{K})$  relationship which is normally desired from any quasi-elastic scattering experiment.

From the calculated quasi-elastic peak shapes generated by the program a calibration curve may be obtained, relating the widths of the calculated peaks  $\Delta E_{calc}$  to the true Lorentzian width  $\Delta E_L$ , for a given incident energy  $E_o$ .

If now the shape of a given experimental quasi-elastic peak, with width  $\Delta E_{obs}$ , is represented accurately by the calculated quasi-elastic peak of identical width, i.e.

$\Delta E_{obs} = \Delta E_{calc}$  then the true Lorentzian width  $\Delta E_L$ , corresponding to  $\Delta E_{obs}$ , may be obtained immediately from inspection of the relevant calibration curve. As the  $\underline{K}$  value of each quasi-elastic peak is known from the particular scattering angle of the relevant detector, the  $\Delta E_L(\underline{K})$  relationship may thus be determined for further analysis.

## 2. Theory

Most theories of quasi-elastic scattering predict a scattering law  $S(\underline{K}, \omega)$  of Lorentzian form [2], [3], [4] independent of the type of sample being analysed. Hence it is assumed here that the experimental quasi-elastic peak shape may be represented by a Lorentzian line shape convoluted by the resolution function of the apparatus.

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(\*) All "widths" described in this report refer to the full width at half peak height.



The line shape of the resolution function is normally taken as the elastic peak of a vanadium time of flight spectrum obtained under identical conditions as the experimental sample spectra. After corrections for the filtering effect have been made the vanadium elastic peaks obtained for different counter angles should be identical since the scattering is incoherent and isotropic. The room temperature Debye-Waller factor of vanadium is given by  $\exp(-0.006 \text{ K}^2)$  and hence varies very little over the normal  $\text{K}^2$  range covered, which is from zero to about  $10 \text{ \AA}^{-2}$ , so that no correction is necessary for this effect.

Alternatively the elastic peak from one of the sample spectra may be used for the resolution function, provided that the spectra were obtained whilst the sample was kept at a temperature which rendered the broadenings of the elastic peaks effectively zero. The difficulty with this method of obtaining the resolution function is that the elastic peaks of the sample spectra may be distorted

- a) from finite coherent scattering of the sample, and
- b) from any remaining inelastic scattering.

Hence the elastic peak shapes at different angles may not be identical, making the choice of the best shape for the resolution function difficult. As vanadium has a negligible coherent scattering cross-section and a region of inelastic scattering, well separated from the elastic peak for normal incident energies of interest, no difficulties arise in this case. However, if the above two effects are not important for the sample of interest, then the alternative method of obtaining the resolution function may be used. This then provides a useful check on the consistency of the apparatus, by comparison with the vanadium results.

With certain chopper-time of flight spectrometers it may be reasonable to represent the resolution function by a Gaussian distribution. In this case the Lorentzian width  $\Delta E_L$  may be obtained directly from tables, (see section 4). In general, however, the resolution function is considerably distorted from a Gaussian shape, depending upon the chopper transmission function and the reactor spectrum, and hence varying according to the incident energy  $E_0$ . The advantage of the present program is that the shape of the resolution function of the instrument can take any form and still allow a  $\Delta E_L$  value to be extracted.

The resolution function is now represented as a function of energy by  $N(E)$ . Folding of this distribution with a Lorentzian energy distribution then gives an energy distribution:

$$T(E) = \int_{-\infty}^{\infty} L(E-E') N(E') dE' \quad [2]$$

where

$$L(E-E') = \frac{\Delta E_L/2}{(\Delta E_L/2)^2 + (E-E')^2} \quad [3]$$

Putting equation [2] in a form suitable for computer calculation we have, as a time of flight distribution

$$T_j(E) = \sum_{i=NS}^{NF} \frac{(\Delta E_L/2) \times N_i(E)}{(\Delta E_L/2)^2 + (E_i - E_j)^2} \quad [4]$$

where  $N_i(E)$  is the value of the resolution function as a time distribution in terms of counts/channel, for the time of flight channel number  $i$  and similarly for  $T_j(E)$  the function produced after  $N_i(E)$  has been convoluted with



the Lorentzian. NS and NF are time of flight channel numbers, chosen symmetrically about  $i_0$ , the channel number corresponding to the incident energy  $E_0$  such that  $(NF-NS) \geq 6\Delta E_R$  where  $\Delta E_R$  is the resolution function width measured in terms of channel numbers. This criterion ensures that the  $T_j(E)$  values obtained from the above integration are sufficiently accurate, for all practical purposes, in representing the complete integration of equation [2].

The time of flight distribution  $T_j(E)$  is calculated and plotted by FOLD for a series of  $\Delta E_L$  values, as a function of channel number  $j$ . In addition  $T_j(E)$  is converted into an energy distribution  $TE_j(E)$  where

$$TE_j(E) = T_j(E) \times j^3 \quad [5]$$

which is also plotted as a function of channel number  $j$ .

Strictly, the resolution function  $N_i(E)$  in equation [4] should first be converted to scattering law form by multiplication with  $(i^4/i_0) \times \exp(-\beta i/2)$

where 
$$\beta i = (E_i - E_0)/kT \quad [6]$$

is the energy transfer for channel  $i$  in units of  $kT$ . The final distribution  $T_j(E)$  will then be in scattering law form which has to be multiplied by  $(i_0/j^4) \times \exp(-\beta i/2)$  for reconversion to a time of flight distribution.

Inclusion of the above factors however, leaves the widths of the functions unchanged and they have therefore been omitted. In addition the  $K$  dependence of the scattering law of the sample, through the Debye-Waller factor  $\exp(-2W)$ , is

omitted in this program. For elastic peaks with wavelength resolutions better than 15% the variation of  $\exp(-2W)$  over the peaks may be considered negligible.

### 3. Program Description.

The program is listed in the Appendix. Most of the variables have comment labels, and will not be discussed further. The data input, with the necessary format is as follows:

CARD 1: (TITLE (I), I = 1,18) (18A4)

This is the title card which is printed out on the first line of the output.

CARDS 2 - 20: (NG(I), I = 1,300) (16I5)

These integer values represent the resolution of the apparatus at the incident energy  $E_0$ , as a function of time of flight channel number  $i$  and correspond to the variable  $N_i(E)$  of equation [4]. Where the resolution function is zero, blanks may be used to fill in the data.

CARD 21: IBAR, CHWD, FPL (I5,2F10.6)

IBAR is the channel number  $i_0$  corresponding to the incident energy  $E_0$ . CHWD is the time analyser channel width in  $\mu\text{sec}$ . FPL is the sample-detector flight path length in metres.

CARD 22: NS, NF (2I5)

These are the channel numbers defining the limits of the numerical integration of equation [4] discussed in section 2.

CARD 23: DELTAE (F10.6)

DELTAE is the width of the Lorentzian function used in equation [4] and is given in eV. As many DELTAE values as required may be fed in on subsequent cards.



CARD 23 + N: 0.2

(F10.6)

N is the total number of cards being used to feed in separate DELTAE values. This final data card acts only as a switch to stop the program.

On output the program prints all the above input parameters together with the incident neutron energy and the corresponding wavelength and reciprocal velocity. Using the plotting routine GRAPH, (5), plots are given of the resolution function and the final Lorentzian broadened function, as both time of flight and energy distributions, between channel numbers NS and NF. The plotted functions are all normalized to unity at maximum intensity and are also printed in a normalized format. The reciprocal velocity in  $\mu\text{sec}/\text{m}$  and energy in meV corresponding to the respective channel numbers are printed at the end of the program for reference.

The widths of the Lorentzian broadened functions are calculated, in channel numbers, by linear interpolation at the half height positions of the peaks, from the time of flight distributions. These widths, together with the corresponding values in meV, are listed in the print out.

Typical FOLD output for a single  $\Delta E_L$  value is shown in the Appendix. This required 0.6 mins for execution.

#### 4. The Calibration Curve.

A few typical Lorentzian broadened resolution curves computed by FOLD are given in Fig. 1 for three different  $\Delta E_L$  values, on a time of flight scale. The input spectrum in this case was chosen to be a perfect Gaussian with an incident energy of 5.24 meV and a resolution of 5.25% in wavelength. The channel width and flight path taken were those currently used on the Double Chopper facility on ISPRA 1 and are 8  $\mu\text{sec}$  and 1.53 m respectively.

Fig. 2 shows a few points, calculated by FOLD, representing the widths  $\Delta E_{\text{calc}}$  of a few typical broadened curves as a function of  $\Delta E_L$  for this Gaussian resolution spectrum. To test the procedure, these points have been compared with a calibration curve, dashed line, obtained from tabulated convolutions of Lorentzian and Gaussian functions [6] using the same Gaussian width as above. Both scales on the graph are given in meV.

Agreement is seen to be good, over the range of broadenings  $\Delta E_L$  shown, for this particular case of a Gaussian input function. Hence it can be taken that FOLD will enable reliable  $\Delta E_L$  values to be obtained, using any arbitrary shaped resolution function.

Fig. 3 illustrates two typical calibration curves for the double chopper facility for two different incident energies of 8.6 meV and 5.2 meV, both with resolutions of 7% in wavelength. Both the Lorentzian and the calculated widths are quoted in meV. The actual resolution functions for these two energies were obtained from vanadium spectra.

Finally it should be noted that the FOLD program may be used quite generally to obtain calibration curves for quasi-elastic scattering from any time of flight apparatus. There is no upper limit on the incident energy which can be used by the program. The lower limit on the incident energy is determined by the number of channels available for the input resolution function. In the present program 300 channels are used but this number can easily be extended if required.

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HENDRICKSON, J.K., WAPD. SR-506 (1954).

Fig.1-LORENTZIAN BROADENED RESOLUTION CURVES  
ON A TIME OF FLIGHT SCALE

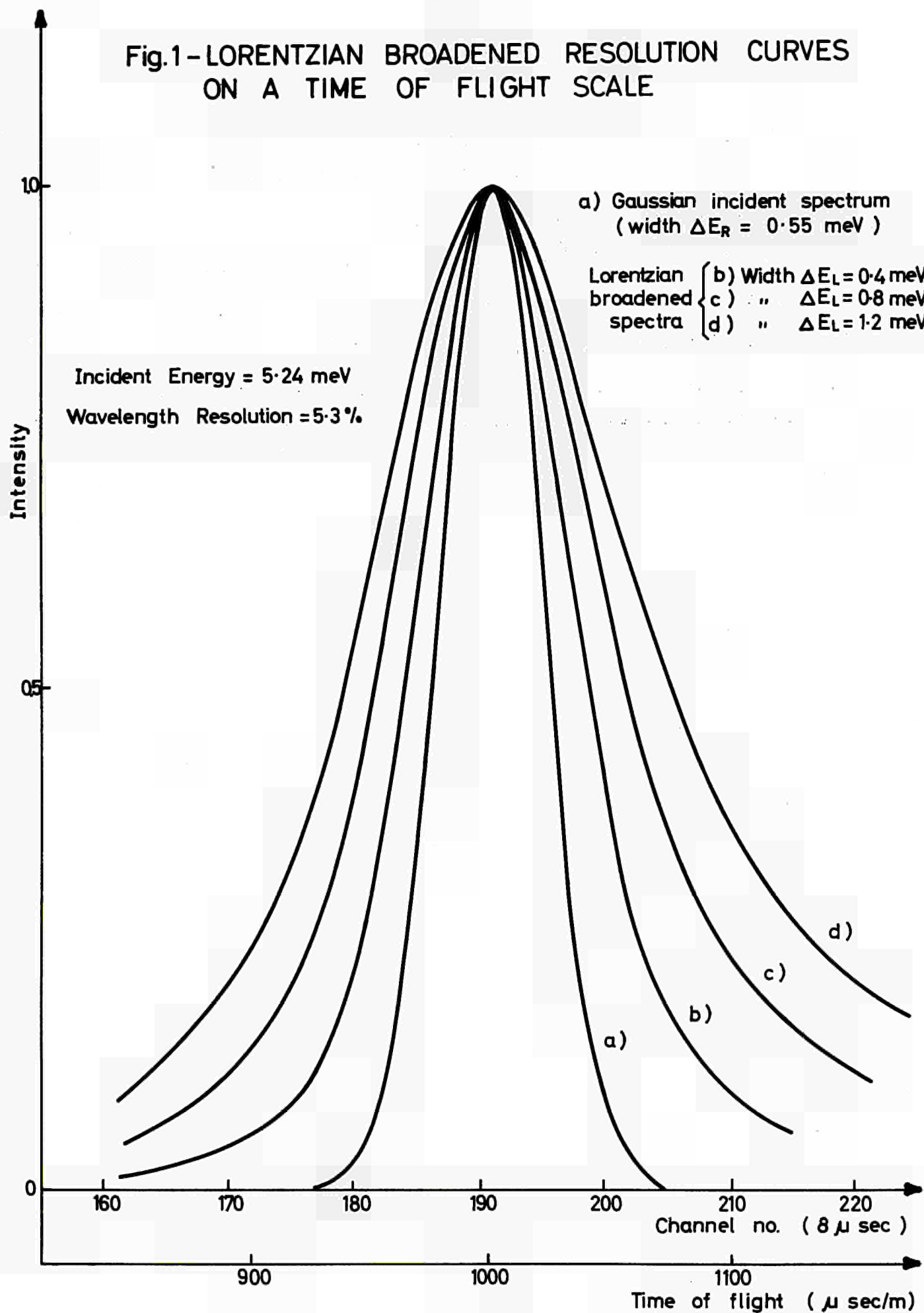




Fig. 2 - COMPARISON OF CALIBRATION OBTAINED FROM FOLD FOR A LORENTZIAN BROADENED GAUSSIAN WITH A TABULATED CALIBRATION CURVE

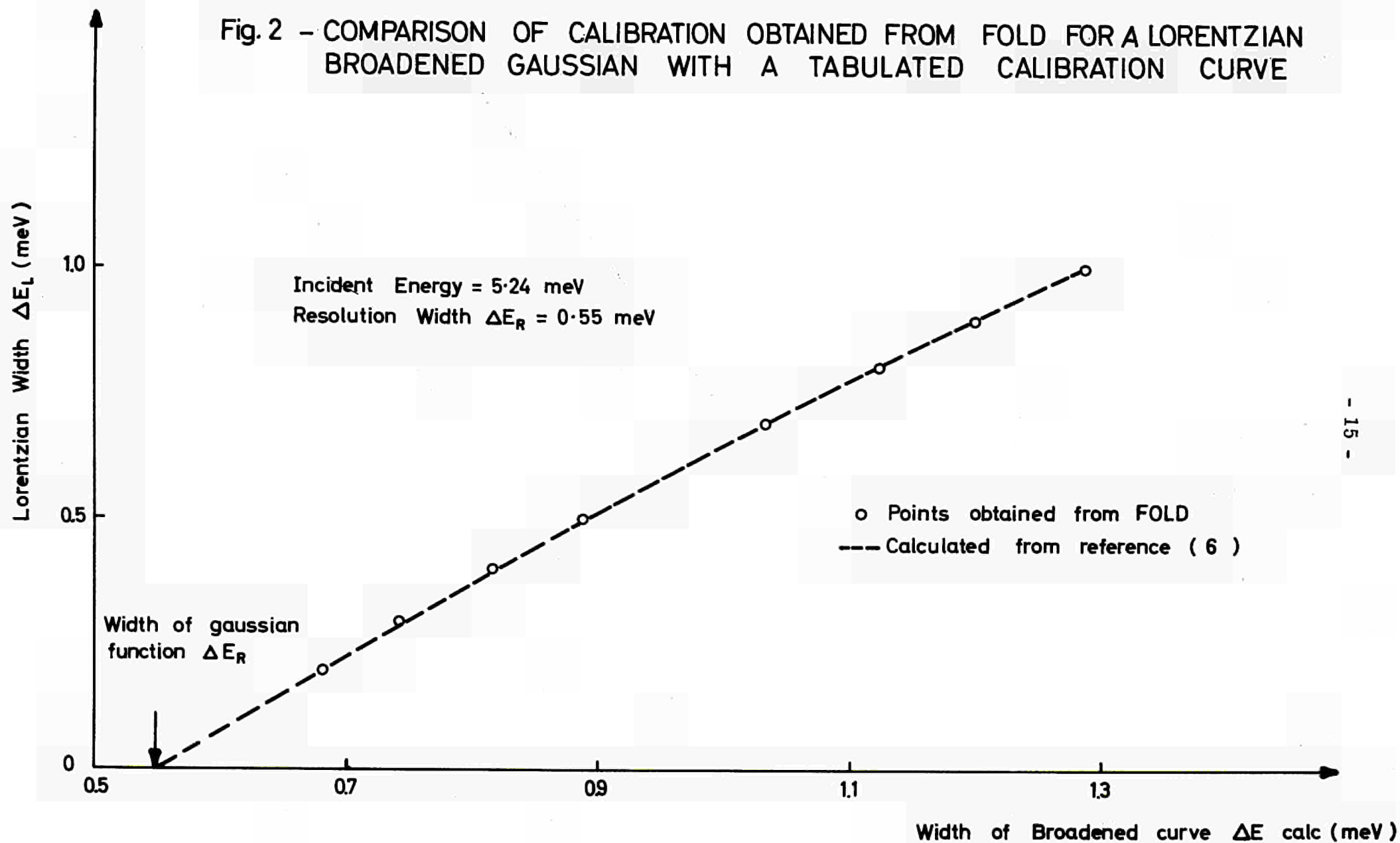
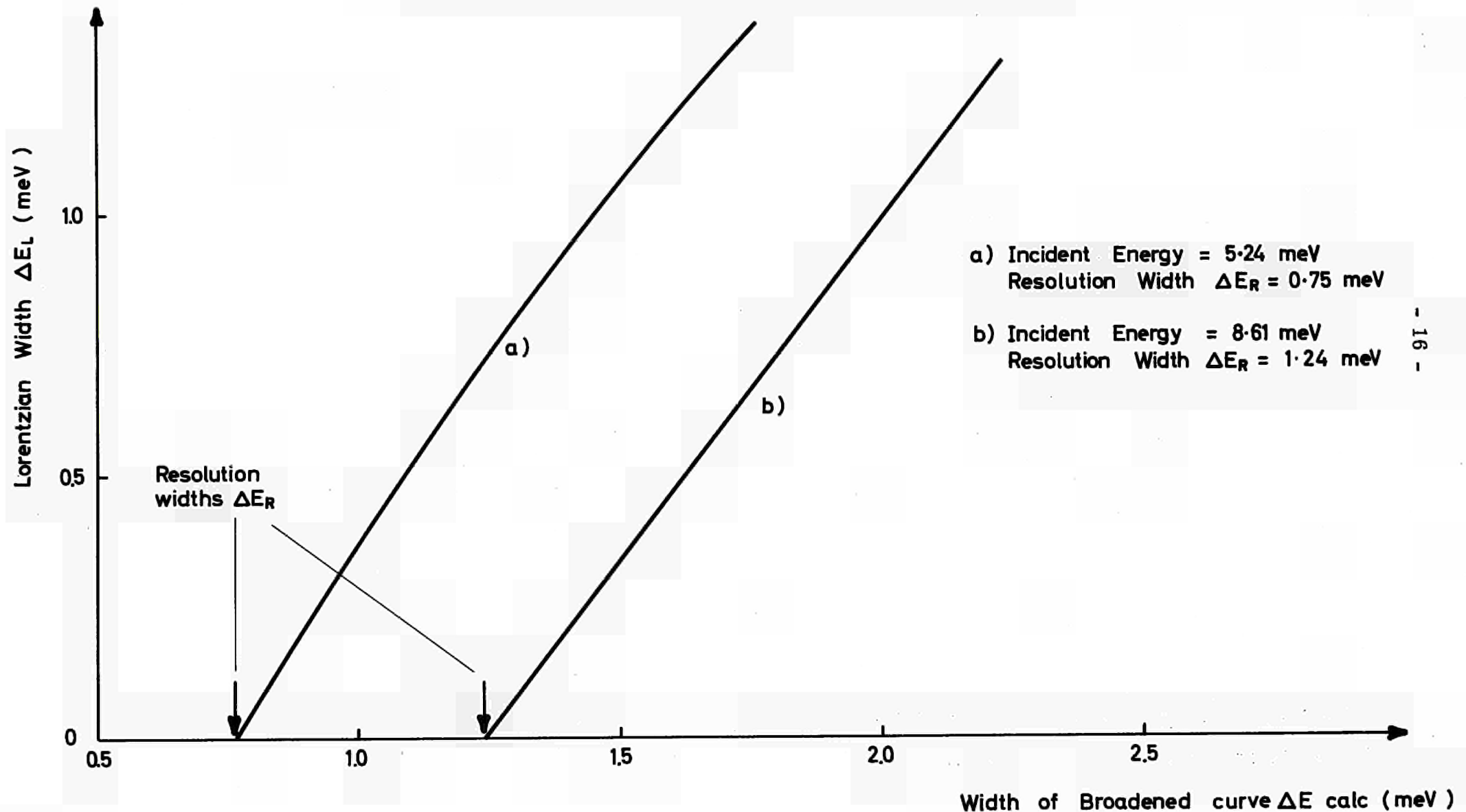


Fig.3 - TYPICAL QUASI ELASTIC CALIBRATION CURVES FOR THE ISPRA 1 DOUBLE CHOPPER FACILITY



APPENDIX

PROGRAM LISTING AND TYPICAL

OUTPUT DATA.

```

      C QUASI ELASTIC NEUTRON SCATTERING CALIBRATION PROGRAM
      C THIS PROGRAM CALCULATES THE BROADENING OF THE INCIDENT SPECTRUM
      C NG(I) WHEN CONVOLUTED WITH A LORENTZIAN SCATTERING LAW S(Q,W) OF
      C ENERGY WIDTH AT HALF HEIGHT OF DELTAE (MEV)
0001      DIMENSION NG(300),S(300),H(300),DEL(300,200),B(300),T(300),E(300),
0002      1,TITL(18),TFHM(300),GE(300),TE(300)
      C CALL INTCTR (180)
0003      READ (5,102) (TITLE(I),I=1,18)
      C READ IN THE INCIDENT SPECTRUM AS A FUNCTION OF CHANNEL NUMBER
0004      READ (5,107) (NG(I),I=1,300)
0005      READ (5,106) IBAR,CHND,FPL
      C IBAR IS THE MEAN CHANNEL NO. OF INCIDENT SPECTRUM
      C CHND IS THE CHANNEL WIDTH (MICROSEC)
      C FPL IS THE FLIGHT PATH LENGTH IN METRES
0006      READ (5,132) IS,IF
0007      READ (5,130) DELTAE
      C DELTAE IS THE FULL WIDTH AT HALF HEIGHT OF THE LORENTZIAN FUNCTION
      C
0008      XBAR=IBAR
0009      B1=10.0*FPL/(CHND*XBAR)
0010      TAU0=10.0/B1
0011      C TAU0 IS THE AVERAGE INCIDENT RECIPROCAL VLOCITY (MICROSEC/METRE)
0012      WAVL0=C.6956/B1
      C WAVL0 IS THE AVERAGE INCIDENT WAVELENGTH (ANGSTROMS)
0013      E1=B1**2
0014      E0=52150.*B1
      C E0 IS THE AVERAGE INCIDENT ENERGY (MEV)
      C
0015      DO 3 I=1,300
0016      XI=I
0017      A=10.0*FPL/(CHND*XI)
0018      TFHM(I)=10.7/A
      C TFHM(I) IS THE RECIPROCAL VELOCITY OF A NEUTRON DETECTED IN CH NO
      C I RELATIVE TO THE INCIDENT CHANNEL NUMBER IBAR
0019      G(I)=NG(I)
0020      GE(I)=G(I)*I**3
      C GE(I) IS THE INCIDENT SPECTRUM ON AN ENERGY SCALE
0021      XMAT=NG(I)
0022      H(I)=XMAT
0023      C 3 CONTINUE
      C
0024      WRITE (6,103) (TITLE(I),I=1,18)
0025      WRITE (6,106) IBAR,CHND,E0,TAU0,WAVL0
0026      WRITE (6,133) IS,IF

```



```

0027      WRITE (6,125)
0028      WRITE (6,136) (S(I),I=1,300)
      C
0029      DO 4 I=1,300
0030      XI=I
0031      XI=XI*0.1*CHLD/FPL
0032      EI=22.25/XI**2
      C
0033      DO 5 J=NS,NF
0034      JJ=J-.45+1
0035      XJ=J
0036      XJ=XJ*0.1*CHLD/FPL
0037      EJ=22.25/XJ**2
      C
0038      DEL(I,JJ)=(EI-EJ)**2
0039      5 CONTINUE
0040      4 CONTINUE
      C
0041      WRITE (6,104) DELTAE
0042      DELTAE=DELTAE/2.0
      C
0043      DO 7 J=1,300
0044      DO 8 K=NS,NF
0045      KK=K-.45+1
0046      XK=K
0047      X=(DELTAE)/(DEL(J,KK)+(DELTAE)**2)
0048      X=X*H(K)
0049      DEL(J,KK)=X
0050      8 CONTINUE
0051      7 CONTINUE
      C
0052      DO 9 S=1,300
0053      S(J)=0.
0054      DO 10 K=NS,NF
0055      KK=K-.45+1
0056      S(J)=S(J)+DEL(J,KK)
0057      10 CONTINUE
0058      TJ=J
0059      TI(J)=S(J)
0060      TI(J)=TI(J)*J**3
      C
0061      TI(J)=15 IS THE BROADENED SPECTRUM ON AN ENERGY SCALE
0062      E(J)=(22.25/XJ**2)
0063      E(J)=E(J)*(10.0*FPL/CHLD)**2
      C
0064      JCL=J.
0065      TIL=0.
0066      JG=0.

```

FORTRAN IV C LEVEL 1, MOD 2

MAIN

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```

0067      Y=C.
0068      C
0069      DO 11 I=1,300
0070      UCL=AMAX1(UCL,GE(I))
0071      TTL=AMAX1(TTL,TE(I))
0072      SC=AMAX1(SC,C(I))
0073      Y=AMAX1(Y,T(I))
0074      11 CONTINUE
0075      C
0076      DO 12 I=1,300
0077      GE(I)=GE(I)/SC
0078      TE(I)=TE(I)/TTE
0079      C(I)=C(I)/SC
0080      T(I)=T(I)/Y
0081      12 CONTINUE
0082      C
0083      WRITE (6,126)
0084      WRITE (6,126) (T(I),I=1,300)
0085      WRITE (6,127)
0086      WRITE (6,126) (TE(I),I=1,300)
0087      C
0088      JP=JF-1.5
0089      C
0090      C
0091      C
0092      C
0093      C
0094      C
0095      C
0096      C
0097      C
0098      C
0099      C
0100      C

```

```

0101      DO 14 I=IBAR,HF
0102      IF (T(I).LT.0.5) GO TO 22
0103      14 CONTINUE
0104      C      22 IMIR=I-1
0105      C      IF ((0.5-T(I)).GT.(T(IMIR)-0.5)) IMIDR=IMIR
0106      C      IF ((0.5-T(I)).LT.(T(IMIR)-0.5)) IMIDR=I
0107      C      IPLUSR=IMIDR+1
0108      C      IMINR=IMIDR-1
0109      C      TANGR=(T(IMINR)-T(IPLUSR))/2.0
0110      C      TMIDR=(T(IMINR)+T(IPLUSR))/2.0
0111      C      FIMINR=IMINR
0112      C      TMIDR=TMIDR-FIMINR
0113      C      WIDTH=TMIDR-TMIDL
0114      C      END OF LINEAR INTERPOLATION ROUTINE
0115      C      EL=52.25/TMIDL**2
0116      C      ER=52.25/TMIDR**2
0117      C      EL=EL*(10.*FPL/CHWD)**2
0118      C      ER=ER*(10.*FPL/CHWD)**2
0119      C      WIDTH=1000.*(EL-ER)
0120      C      DELTAE=DELTA*2000.
0121      C      WRITE (6,140) WIDTH,DELTAE
0122      C      WIDTH IS WIDTH OF BROADENED PEAK IN MEV
0123      C      WRITE (6,141) WIDTH
0124      C      WIDTH IS WIDTH OF BROADENED PEAK IN CHANNEL NUMBERS
0125      C      READ IN NEXT LORENTZIAN WIDTH DELTAE (EV)
0126      C      READ (5,130) DELTAE
0127      C      IF (DELTAE-2.) 6,23,23
0128      C      23 WRITE (6,142)
0129      C      DO 15 J=1,200
0130      C      WRITE (6,143) (J,TF.LH(J),E(J))
0131      C      15 CONTINUE
0132      C      101 FORMAT (18A4)
0133      C      102 FORMAT (1H1,13A4)
0134      C      103 FORMAT (1H1,18X,50H F.W.H.H. OF THE LORENTZIAN FUNCTION IS DELTAE
0135      C      1 = ,F4.1,3X,10H(EV))
0136      C      104 FORMAT (13,2F10.6)
0137      C      105 FORMAT (13,15)
0138      C      106 FORMAT (77/10X,24H INCIDENT CHANNEL NUMBER=,17,2X,14H CHANNEL WIDTH=
0139      C      1,F4.1,1X,30H MICROSEC,3X,16H INCIDENT ENERGY=,F7.4,2X,30H MEV,77/10X,15H

```

```

      TIME OF FLIGHT=,2X,F7.2,2X,10MICROSEC/M,5X,11HWAVELENGTH=,F8.4,2X
      3,911A,10STRJMS)
0134      109 FORMAT (///30H WIDTH OF THE BROADENED FUNCTION IS ,3X,F10.3,3X,15H
      10 CHANNEL NUMBERS///)
0135      110 FORMAT (///39H ENERGY WIDTH OF BROADENED FUNCTION IS ,F10.6,3X,3HM
      1EV,3X,34H CORRESPONDING LORENTZIAN WIDTH WAS,3X,F10.6,3X,3HMEV///)
0136      125 FORMAT (///24X,18H INCIDENT SPECTRUM,/)
0137      129 FORMAT (10(3X,F7.4))
0138      127 FORMAT (///10X,27H ENERGY SPECTRUM...BROADENED//)
0139      128 FORMAT (///10X,35H TIME OF FLIGHT SPECTRUM...BROADENED//)
0140      150 FORMAT (F10.6)
0141      152 FORMAT (2I5)
0142      156 FORMAT (10(2X,F8.2))
0143      158 FORMAT (710X,40H TRUNCATION TAKEN BETWEEN CHANNEL NUMBERS,17,2X,3HA
      1ND,17//)
0144      140 FORMAT (///10X,62H PLOT OF INCIDENT AND BROADENED SPECTRA ON TIME O
      1F FLIGHT SCALE,72X,32H (ABSCISSAE IS IN MICROSEC/METRE)//)
0145      141 FORMAT (///10X,57H PLOT OF INCIDENT AND BROADENED SPECTRA ON AN ENE
      1RGY SCALE,77X,21H (ABSCISSAE IS IN EV)//)
0146      142 FORMAT (11H,15X,11H CHANNEL NO.,10X,27H TIME OF FLIGHT (MICROSEC/M),
      11X,11H ENERGY (EV) //)
0147      143 FORMAT (13X,15,19X,F10.3,19X,F13.6)
      C
0148      . STOP
0149      . END

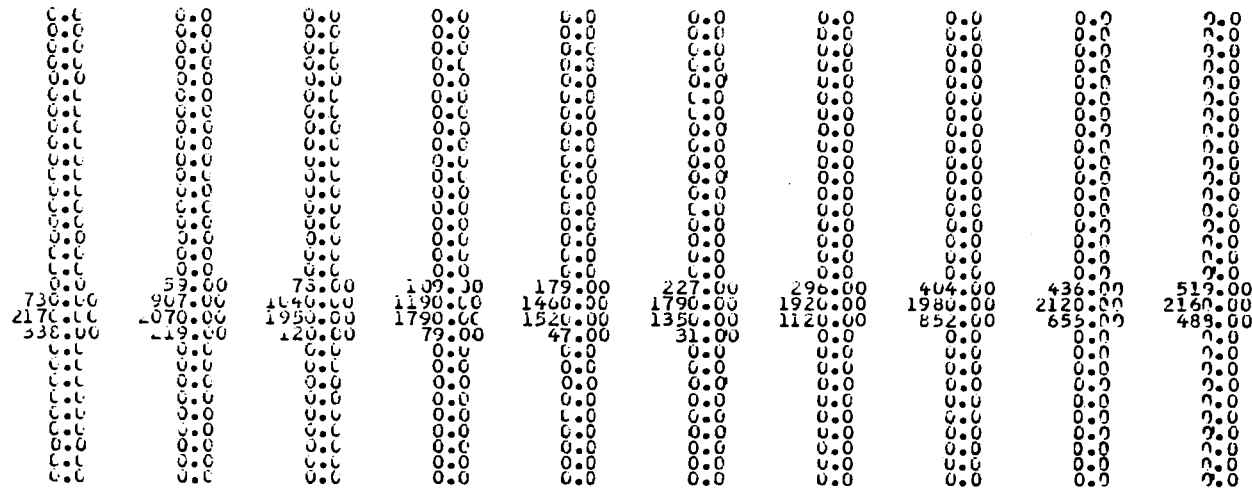
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QUASI ELASTIC SCATTERING CALIBRATION FOR DOUBLE CHOPPER DATA

INCIDENT CHANNEL NUMBER= 191 CHANNEL WIDTH= 8.0 MICROSEC INCIDENT ENERGY= 5.2367 MEV  
 TIME OF FLIGHT= 998.69 MICROSEC/H WAVELENGTH= 3.9528 ANGSTROMS  
 TRUNCATION TAKEN BETWEEN CHANNEL NUMBERS 151 AND 231

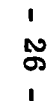
INCIDENT SPECTRUM



- 24 -

[illegible][illegible]

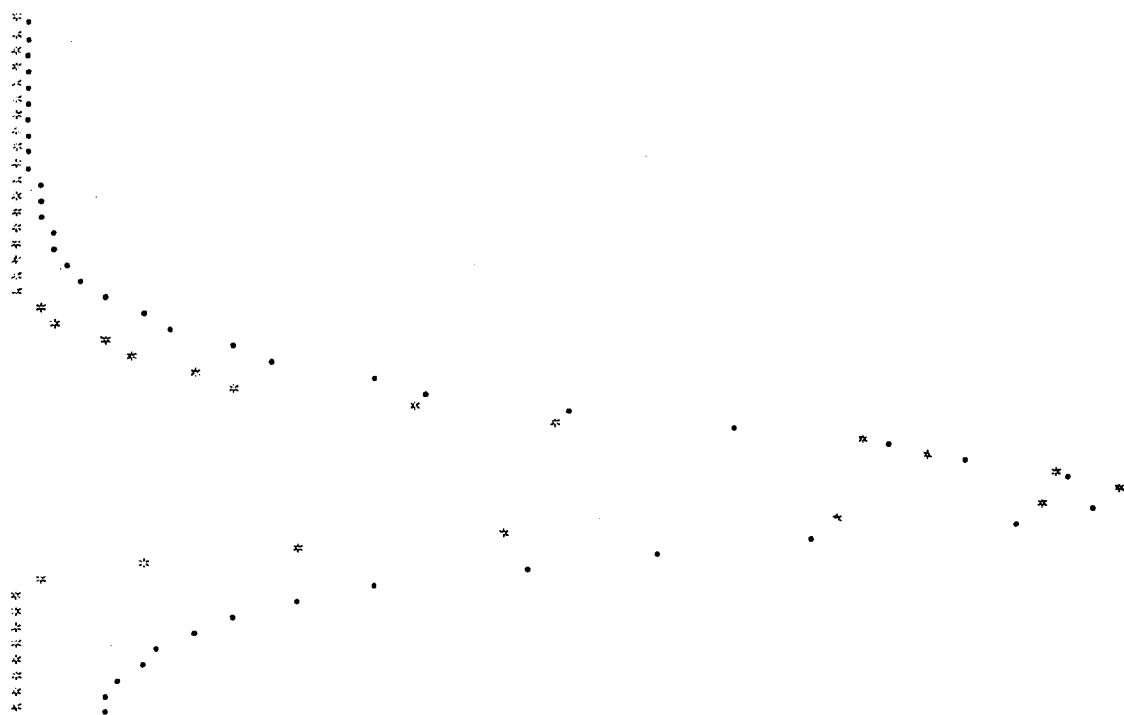
0.4377	0.5044	0.5750	0.6494	0.7237	0.7949	0.8596	0.9146	0.9575	0.9865
0.4400	0.5070	0.5775	0.6519	0.7262	0.7974	0.8621	0.9171	0.9600	0.9890
0.4423	0.5093	0.5800	0.6542	0.7285	0.8000	0.8647	0.9194	0.9623	0.9913
0.4446	0.5116	0.5825	0.6565	0.7308	0.8023	0.8670	0.9217	0.9646	0.9936
0.4469	0.5139	0.5850	0.6588	0.7331	0.8046	0.8693	0.9240	0.9669	0.9959
0.4492	0.5162	0.5875	0.6611	0.7354	0.8069	0.8716	0.9263	0.9692	0.9982
0.4515	0.5185	0.5900	0.6634	0.7377	0.8092	0.8739	0.9286	0.9715	1.0005
0.4538	0.5208	0.5925	0.6657	0.7400	0.8115	0.8762	0.9309	0.9738	1.0028
0.4561	0.5231	0.5950	0.6680	0.7423	0.8138	0.8785	0.9332	0.9761	1.0051
0.4584	0.5254	0.5975	0.6703	0.7446	0.8161	0.8808	0.9355	0.9784	1.0074
0.4607	0.5277	0.6000	0.6726	0.7469	0.8184	0.8831	0.9378	0.9807	1.0097
0.4630	0.5300	0.6025	0.6749	0.7492	0.8207	0.8854	0.9401	0.9830	1.0120
0.4653	0.5323	0.6050	0.6772	0.7515	0.8230	0.8877	0.9424	0.9853	1.0143
0.4676	0.5346	0.6075	0.6795	0.7538	0.8253	0.8900	0.9447	0.9876	1.0166
0.4699	0.5369	0.6100	0.6818	0.7561	0.8276	0.8923	0.9470	0.9899	1.0189
0.4722	0.5392	0.6125	0.6841	0.7584	0.8299	0.8946	0.9493	0.9922	1.0212
0.4745	0.5415	0.6150	0.6864	0.7607	0.8322	0.8969	0.9516	0.9945	1.0235
0.4768	0.5438	0.6175	0.6887	0.7630	0.8345	0.8992	0.9539	0.9968	1.0258
0.4791	0.5461	0.6200	0.6910	0.7653	0.8368	0.9015	0.9562	0.9991	1.0281
0.4814	0.5484	0.6225	0.6933	0.7676	0.8391	0.9038	0.9585	1.0014	1.0304
0.4837	0.5507	0.6250	0.6956	0.7699	0.8414	0.9061	0.9608	1.0037	1.0327
0.4860	0.5530	0.6275	0.6979	0.7722	0.8437	0.9084	0.9631	1.0060	1.0350
0.4883	0.5553	0.6300	0.7002	0.7745	0.8460	0.9107	0.9654	1.0083	1.0373
0.4906	0.5576	0.6325	0.7025	0.7768	0.8483	0.9130	0.9677	1.0106	1.0396
0.4929	0.5599	0.6350	0.7048	0.7791	0.8506	0.9153	0.9700	1.0129	1.0419
0.4952	0.5622	0.6375	0.7071	0.7814	0.8529	0.9176	0.9723	1.0152	1.0442
0.4975	0.5645	0.6400	0.7094	0.7837	0.8552	0.9199	0.9746	1.0175	1.0465
0.4998	0.5668	0.6425	0.7117	0.7860	0.8575	0.9222	0.9769	1.0198	1.0488
0.5021	0.5691	0.6450	0.7140	0.7883	0.8598	0.9245	0.9792	1.0221	1.0511
0.5044	0.5714	0.6475	0.7163	0.7906	0.8621	0.9268	0.9815	1.0244	1.0534
0.5067	0.5737	0.6500	0.7186	0.7929	0.8644	0.9291	0.9838	1.0267	1.0557
0.5090	0.5760	0.6525	0.7209	0.7952	0.8667	0.9314	0.9861	1.0290	1.0580
0.5113	0.5783	0.6550	0.7232	0.7975	0.8690	0.9337	0.9884	1.0313	1.0603
0.5136	0.5806	0.6575	0.7255	0.7998	0.8713	0.9360	0.9907	1.0336	1.0626
0.5159	0.5829	0.6600	0.7278	0.8021	0.8736	0.9383	0.9930	1.0359	1.0649
0.5182	0.5852	0.6625	0.7301	0.8044	0.8759	0.9406	0.9953	1.0382	1.0672
0.5205	0.5875	0.6650	0.7324	0.8067	0.8782	0.9429	0.9976	1.0405	1.0695
0.5228	0.5898	0.6675	0.7347	0.8090	0.8805	0.9452	0.9999	1.0428	1.0718
0.5251	0.5921	0.6700	0.7370	0.8113	0.8828	0.9475	1.0022	1.0451	1.0741
0.5274	0.5944	0.6725	0.7393	0.8136	0.8851	0.9498	1.0045	1.0474	1.0764
0.5297	0.5967	0.6750	0.7416	0.8159	0.8874	0.9521	1.0068	1.0497	1.0787
0.5320	0.5990	0.6775	0.7439	0.8182	0.8897	0.9544	1.0091	1.0520	1.0810
0.5343	0.6013	0.6800	0.7462	0.8205	0.8920	0.9567	1.0114	1.0543	1.0833
0.5366	0.6036	0.6825	0.7485	0.8228	0.8943	0.9590	1.0137	1.0566	1.0856
0.5389	0.6059	0.6850	0.7508	0.8251	0.8966	0.9613	1.0160	1.0589	1.0879
0.5412	0.6082	0.6875	0.7531	0.8274	0.8989	0.9636	1.0183	1.0612	1.0902
0.5435	0.6105	0.6900	0.7554	0.8297	0.9012	0.9659	1.0206	1.0635	1.0925
0.5458	0.6128	0.6925	0.7577	0.8320	0.9035	0.9682	1.0229	1.0658	1.0948
0.5481	0.6151	0.6950	0.7600	0.8343	0.9058	0.9705	1.0252	1.0681	1.0971
0.5504	0.6174	0.6975	0.7623	0.8366	0.9081	0.9728	1.0275	1.0704	1.0994
0.5527	0.6197	0.7000	0.7646	0.8389	0.9104	0.9751	1.0298	1.0727	1.1017
0.5550	0.6220	0.7025	0.7669	0.8412	0.9127	0.9774	1.0321	1.0750	1.1040
0.5573	0.6243	0.7050	0.7692	0.8435	0.9150	0.9797	1.0344	1.0773	1.1063
0.5596	0.6266	0.7075	0.7715	0.8458	0.9173	0.9820	1.0367	1.0796	1.1086
0.5619	0.6289	0.7100	0.7738	0.8481	0.9196	0.9843	1.0390	1.0819	1.1109
0.5642	0.6312	0.7125	0.7761	0.8504	0.9219	0.9866	1.0413	1.0842	1.1132
0.5665	0.6335	0.7150	0.7784	0.8527	0.9242	0.9889	1.0436	1.0865	1.1155
0.5688	0.6358	0.7175	0.7807	0.8550	0.9265	0.9912	1.0459	1.0888	1.1178
0.5711	0.6381	0.7200	0.7830	0.8573	0.9288	0.9935	1.0482	1.0911	1.1201
0.5734	0.6404	0.7225	0.7853	0.8596	0.9311	0.9958	1.0505	1.0934	1.1224
0.5757	0.6427	0.7250	0.7876	0.8619	0.9334	0.9981	1.0528	1.0957	1.1247
0.5780	0.6450	0.7275	0.7899	0.8642	0.9357	1.0004	1.0551	1.0980	1.1270
0.5803	0.6473	0.7300	0.7922	0.8665	0.9380	1.0027	1.0574	1.1003	1.1293
0.5826	0.6496	0.7325	0.7945	0.8688	0.9403	1.0050	1.0597	1.1026	1.1316
0.5849	0.6519	0.7350	0.7968	0.8711	0.9426	1.0073	1.0620	1.1049	1.1339
0.5872	0.6542	0.7375	0.7991	0.8734	0.9449	1.0096	1.0643	1.1072	1.1362
0.5895	0.6565	0.7400	0.8014	0.8757	0.9472	1.0119	1.0666	1.1095	1.1385
0.5918	0.6588	0.7425	0.8037	0.8780	0.9495	1.0142	1.0689	1.1118	1.1408
0.5941	0.6611	0.7450	0.8060	0.8803	0.9518	1.0165	1.0712	1.1141	1.1431
0.5964	0.6634	0.7475	0.8083	0.8826	0.9541	1.0188	1.0735	1.1164	1.1454
0.5987	0.6657	0.7500	0.8106	0.8849	0.9564	1.0211	1.0758	1.1187	1.1477
0.6010	0.6680	0.7525	0.8129	0.8872	0.9587	1.0234	1.0781	1.1210	1.1500
0.6033	0.6703	0.7550	0.8152	0.8895	0.9610	1.0257	1.0804	1.1233	1.1523
0.6056	0.6726	0.7575	0.8175	0.8918	0.9633	1.0280	1.0827	1.1256	1.1546
0.6079	0.6749	0.7600	0.8198	0.8941	0.9656	1.0303	1.0850	1.1279	1.1569
0.6102	0.6772	0.7625	0.8221	0.8964	0.9679	1.0326	1.0873	1.1302	1.1592
0.6125	0.6795	0.7650	0.8244	0.8987	0.9702	1.0349	1.0896	1.1325	1.1615
0.6148	0.6818	0.7675	0.8267	0.9010	0.9725	1.0372	1.0919	1.1348	1.1638
0.6171	0.6841	0.7700	0.8290	0.9033	0.9748	1.0395	1.0942	1.1371	1.1661
0.6194	0.6864	0.7725	0.8313	0.9056	0.9771	1.0418	1.0965	1.1394	1.1684
0.6217	0.6887	0.7750	0.8336	0.9079	0.9794	1.0441	1.0988	1.1417	1.1707
0.6240	0.6910	0.7775	0.8359	0.9102	0.9817	1.0464	1.1011	1.1440	1.1730
0.6263	0.6933	0.7800	0.8382	0.9125	0.9840	1.0487	1.1034	1.1463	1.1753
0.6286	0.6956	0.7825	0.8405	0.9148	0.9863	1.0510	1.1057	1.1486	1.1776
0.6309	0.6979	0.7850	0.8428	0.9171	0.9886	1.0533	1.1080	1.1509	1.1799
0.6332	0.7002	0.7875	0.8451	0.9194	0.9909	1.0556	1.1103	1.1532	1.1822
0.6355	0.7025	0.7900	0.8474	0.9217	0.9932	1.0579	1.1126	1.1555	1.1845
0.6378	0.7048	0.7925	0.8497	0.9240	0.9955	1.0602	1.1149	1.1578	1.1868
0.6401	0.7071	0.7950	0.8520	0.9263	0.9978	1.0625	1.1172	1.1601	1.1891
0.6424	0.7094	0.7975	0.8543	0.9286	1.0001	1.0648	1.1195	1.1624	1.1914
0.6447	0.7117	0.8000	0.8566	0.9309	1.0024	1.0671	1.1218	1.1647	1.1937
0.6470	0.7140	0.8025	0.8589	0.9332	1.0047	1.0694	1.1241	1.1670	1.1960
0.6493	0.7163	0.8050	0.8612	0.9355	1.0070	1.0717	1.1264	1.1693	1.1983
0.6516	0.7186	0.8075	0.8635	0.9378	1.0093	1.0740	1.1287	1.1716	1.2006
0.6539	0.7209	0.8100	0.8658	0.9401	1.0116	1.0763	1.1310	1.1739	1.2029
0.6562	0.7232	0.8125	0.8681	0.9424	1.0139	1.0786	1.1333	1.1762	1.2052
0.6585	0.7255	0.8150	0.8704	0.9447	1.0162	1.0809	1.1356	1.1785	1.2075
0.6608	0.7278	0.8175	0.8727	0.9470	1.0185	1.0832	1.1379	1.1808	1.2098
0.6631	0.7301	0.8200	0.8750	0.9493	1.0208	1.0855	1.1402	1.1831	1.2121
0.6654	0.7324	0.8225	0.8773	0.9516	1.0231	1.0878	1.1425	1.1854	1.2144
0.6677	0.7347	0.8250	0.8796	0.9539	1.0254	1.0901	1.1448	1.1877	1.2167
0.6700	0.7370	0.8275	0.8819	0.9562	1.0277	1.0924	1.1471	1.1900	1.2190
0.6723	0.7393	0.8300	0.8842	0.9585	1.0300	1.0947	1.1494	1.1923	1.2213
0.6746	0.7416	0.8325	0.8865	0.9608	1.0323	1.0970	1.1517	1.1946	1.2236
0.6769	0.7439	0.8350	0.8888	0.9631	1.0346	1.0993	1.1540	1.1969	1.2259
0.6792	0.7462	0.8375	0.8911	0.9654	1.0369	1.1016	1.1563	1.1992	1.2282
0.6815	0.7485	0.8400	0.8934	0.9677	1.0392	1.1039	1.1586	1.2015	1.2305
0.6838	0.7508	0.8425	0.8957						



INFORMATION ON ORDINATE SCALE - LEFT SCALING POINT = 0.0 =INTERVAL BETWEEN SCALING POINTS= 1.000E-01



3.50E-01	3.42E-01	0.1
3.47E-01	3.41E-01	0.0
3.45E-01	3.40E-01	0.0
3.43E-01	3.39E-01	0.0
3.41E-01	3.38E-01	0.0
3.39E-01	3.37E-01	0.0
3.37E-01	3.36E-01	0.0
3.35E-01	3.35E-01	0.0
3.33E-01	3.34E-01	0.0
3.31E-01	3.33E-01	0.0
3.29E-01	3.32E-01	0.0
3.27E-01	3.31E-01	0.0
3.25E-01	3.30E-01	0.0
3.23E-01	3.29E-01	0.0
3.21E-01	3.28E-01	0.0
3.19E-01	3.27E-01	0.0
3.17E-01	3.26E-01	0.0
3.15E-01	3.25E-01	0.0
3.13E-01	3.24E-01	0.0
3.11E-01	3.23E-01	0.0
3.09E-01	3.22E-01	0.0
3.07E-01	3.21E-01	0.0
3.05E-01	3.20E-01	0.0
3.03E-01	3.19E-01	0.0
3.01E-01	3.18E-01	0.0
2.99E-01	3.17E-01	0.0
2.97E-01	3.16E-01	0.0
2.95E-01	3.15E-01	0.0
2.93E-01	3.14E-01	0.0
2.91E-01	3.13E-01	0.0
2.89E-01	3.12E-01	0.0
2.87E-01	3.11E-01	0.0
2.85E-01	3.10E-01	0.0
2.83E-01	3.09E-01	0.0
2.81E-01	3.08E-01	0.0
2.79E-01	3.07E-01	0.0
2.77E-01	3.06E-01	0.0
2.75E-01	3.05E-01	0.0
2.73E-01	3.04E-01	0.0
2.71E-01	3.03E-01	0.0
2.69E-01	3.02E-01	0.0
2.67E-01	3.01E-01	0.0
2.65E-01	3.00E-01	0.0
2.63E-01	2.99E-01	0.0
2.61E-01	2.98E-01	0.0
2.59E-01	2.97E-01	0.0
2.57E-01	2.96E-01	0.0
2.55E-01	2.95E-01	0.0
2.53E-01	2.94E-01	0.0
2.51E-01	2.93E-01	0.0
2.49E-01	2.92E-01	0.0
2.47E-01	2.91E-01	0.0
2.45E-01	2.90E-01	0.0
2.43E-01	2.89E-01	0.0
2.41E-01	2.88E-01	0.0
2.39E-01	2.87E-01	0.0
2.37E-01	2.86E-01	0.0
2.35E-01	2.85E-01	0.0
2.33E-01	2.84E-01	0.0
2.31E-01	2.83E-01	0.0
2.29E-01	2.82E-01	0.0
2.27E-01	2.81E-01	0.0
2.25E-01	2.80E-01	0.0
2.23E-01	2.79E-01	0.0
2.21E-01	2.78E-01	0.0
2.19E-01	2.77E-01	0.0
2.17E-01	2.76E-01	0.0
2.15E-01	2.75E-01	0.0
2.13E-01	2.74E-01	0.0
2.11E-01	2.73E-01	0.0
2.09E-01	2.72E-01	0.0
2.07E-01	2.71E-01	0.0
2.05E-01	2.70E-01	0.0
2.03E-01	2.69E-01	0.0
2.01E-01	2.68E-01	0.0
1.99E-01	2.67E-01	0.0
1.97E-01	2.66E-01	0.0
1.95E-01	2.65E-01	0.0
1.93E-01	2.64E-01	0.0
1.91E-01	2.63E-01	0.0
1.89E-01	2.62E-01	0.0
1.87E-01	2.61E-01	0.0
1.85E-01	2.60E-01	0.0
1.83E-01	2.59E-01	0.0
1.81E-01	2.58E-01	0.0
1.79E-01	2.57E-01	0.0
1.77E-01	2.56E-01	0.0
1.75E-01	2.55E-01	0.0
1.73E-01	2.54E-01	0.0
1.71E-01	2.53E-01	0.0
1.69E-01	2.52E-01	0.0
1.67E-01	2.51E-01	0.0
1.65E-01	2.50E-01	0.0
1.63E-01	2.49E-01	0.0
1.61E-01	2.48E-01	0.0
1.59E-01	2.47E-01	0.0
1.57E-01	2.46E-01	0.0
1.55E-01	2.45E-01	0.0
1.53E-01	2.44E-01	0.0
1.51E-01	2.43E-01	0.0
1.49E-01	2.42E-01	0.0
1.47E-01	2.41E-01	0.0
1.45E-01	2.40E-01	0.0
1.43E-01	2.39E-01	0.0
1.41E-01	2.38E-01	0.0
1.39E-01	2.37E-01	0.0
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ENERGY WIDTH OF BROADENED FUNCTION IS 1.079267 MEV CORRESPONDING LORENTZIAN WIDTH HAS 0.500000 MEV

WIDTH OF THE BROADENED FUNCTION IS 19.421 CHANNEL NUMBERS

CHANNEL NO.

TIME OF FLIGHT (MICROSEC/NI)

ENERGY (EVI)

1	5.229	191.039230
2	10.458	47.759790
3	15.686	21.226578
4	20.915	11.939952
5	26.144	7.641566
6	31.373	5.306643
7	36.601	3.898759
8	41.830	2.984987
9	47.059	2.358508
10	52.288	1.916392
11	57.516	1.578836
12	62.745	1.326661
13	67.974	1.130409
14	73.203	0.974690
15	78.431	0.849063
16	83.660	0.746247
17	88.889	0.661035
18	94.118	0.589627
19	99.346	0.529194
20	104.575	0.477598
21	109.804	0.433195
22	115.033	0.394709
23	120.262	0.361133
24	125.490	0.331665
25	130.719	0.305663
26	135.948	0.282602
27	141.177	0.262056
28	146.405	0.243672
29	151.634	0.227157
30	156.863	0.212266
31	162.092	0.198792
32	167.320	0.186562
33	172.549	0.175426
34	177.778	0.165259
35	183.007	0.155950
36	188.235	0.147407
37	193.464	0.139547
38	198.693	0.132299
39	203.922	0.125601
40	209.150	0.119399
41	214.379	0.113646
42	219.608	0.108299
43	224.837	0.103320
44	230.065	0.098677
45	235.294	0.094340
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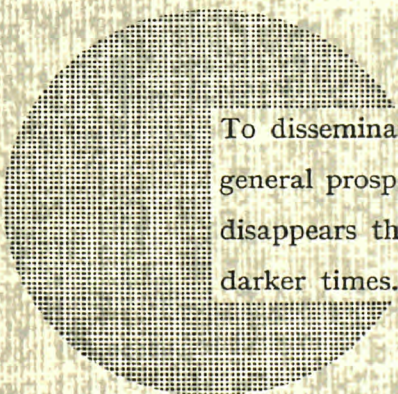
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Alfred Nobel



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